

# In Favour of a Schrödinger Evolution for the Universe

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# Fundamental Postulate

It is possible, in principle, to measure time using a clock.

$\Rightarrow$  clock readings are observables.



# Negative Argument for Reduced Phase Space and Dirac Quantization

Clock readings are observable



Global reparametrizations do not reflect a redundancy at the level of phase space variables.

## Space of Solutions

What the space of solutions is:

- The space of valid initial data for the classical evolution problem.  
⇒ meaningless without a notion of evolution.

AND NOT

- A physical (i.e., reduced) phase space (no evolution).  
⇒ consequence of a redundancy.

Thus:

A quantization of the space of solutions does not capture the physics of our Fundamental Postulate ⇒ clocks not treated as observable.

This applies to:

- Reduced phase space quantization.
- Dirac quantization.



# Our Proposal (In Brief)

For globally reparametrization invariant theories:

## Relational Quantization

- Observables (mutables): Full non-redundant phase space,  $(q, p)$ .
- Classical evolution: integral curves of the Hamiltonian constraint,  $\mathcal{H}$ , labelled by unobservable parameter  $\tau$

$$\dot{q} = \{q, \mathcal{H}\} \quad \dot{p} = \{p, \mathcal{H}\}$$

- Quantum evolution: mutable algebra evolves according to

$$\hat{q}(\tau) = e^{-i\tau\hat{\mathcal{H}}}\hat{q}(0)e^{i\tau\hat{\mathcal{H}}} \quad \hat{p}(\tau) = e^{-i\tau\hat{\mathcal{H}}}\hat{p}(0)e^{i\tau\hat{\mathcal{H}}}$$

⇒ Time Dependent Schrödinger equation.

For gravity:

## Shape Dynamics

- Fix local time by requiring local scale invariance ⇒ Shape Dynamics.
- Relationally quantize Shape Dynamics.



# The “Global” in Global Reparametrization

## Invariance of Action

Action  $S = \int d\tau L(q, \dot{q})$  invariant under

$$\tau \rightarrow f(\tau)$$

Note:

- $\tau$  - Single time variable (not a function of  $\Sigma$ ).
- $q$  - can be a field on  $\Sigma$ .

## Caveat

$L$  - homogeneous of degree 1 in  $\dot{q}$ .

$\Rightarrow$  Hamiltonian constraint  $\mathcal{H} = 0$ .

Classical solutions are integral curves of  $\mathcal{H}$  labelled by  $\tau$ .



# Reduced Phase Space Quantization

Note: will distinguish from deparametrization.

## Observation

Classical histories = integral curves of  $\mathcal{H}$



- Quotienting flow of  $\mathcal{H} \Rightarrow$  space of initial data.
- No Hamiltonian flow on reduced phase space.

$\Rightarrow$  cannot generate classical solutions. (Clocks are not redundancies.)

$\therefore$  inappropriate for quantization. (Only appropriate for a genuine redundancy.)

Space of initial data is isomorphic to, but not have the same representational capacity as, the space of histories.



# Deparametrization Definition

Note: will consider distinct from phase space reduction.

## Nomenclature

Partition phase space into:

- potential clocks (partial observables).
- Dirac observables of a particular clock (complete observables).

Procedure: (Donald's talk)

- Choose a clock  $T \equiv q_i \in q$
- Identify  $T = \tau$  (special gauge fixing of  $\mathcal{H} = 0$ )
- Find Dirac observables  $Q_I = q_I, P_I = p_i$  for  $(I \neq i)$ .
- The  $(Q, P)$  are the complete observables with respect to the partial observable  $T$ .
- Solve  $\mathcal{H} = 0$  for  $E \equiv p_i = H_i$ .
- Use  $H_i$  to evolve  $(Q, P)$ .

$\Rightarrow$  Gives (correct) classical evolution of  $(Q, P)$  in terms of  $T$ .



# Deparametrization Difficulties

Leads to effective evolution.

## But...

- Good clocks are hard (/impossible) to find. (E.g., Philipp's talk)
- Dirac observables don't tell the whole story.  
⇒  $T$  is not a Dirac observable. (Spacetime scalars??)
- $T$  has an ambiguous interpretation (Rovelli: observable but not predictable).
- Wrong physical picture for Yang–Mills (phase space reduction is).

Difficulties are interpretational in classical theory, but physical in quantum.





# Quantum Deparametrization Definition

Choose a particular deparametrization wrt  $T$ :

- Complete (Dirac) observables:

$$(Q, P); \{Q, P\} = \mathbb{1} \quad \Rightarrow \quad (\hat{Q}, \hat{P}); [\hat{Q}, \hat{P}] = i\hbar \hat{\mathbb{1}}$$

- Evolution:

$$E = H_T(Q, P, T) \quad \Rightarrow \quad -i\hbar \frac{\partial \psi}{\partial T} = \hat{H}_T(\hat{Q}, \hat{P}, T)$$

$\Rightarrow$  gives evolution of  $(\hat{Q}, \hat{P})$  in terms of clock parameter  $T$ .



# Quantum Deparametrization Difficulties

Leads to effective notion of evolution.

But...

- Good clocks are hard (/impossible) to find.  
⇒ Serious (/insurmountable) ordering ambiguities.
- Different clock choices are unitarily inequivalent.  
⇒  $T$  takes classical values, but is an operator for different clock choice.
- Conflict: Dirac observables violate our Fundamental Principle.
- The status of  $T$  is physically restrictive in quantum theory.

⇒ Need a way to represent all phase space functions as quantum observables  
AND recover the correct classical limit.



# Classical Ontology

## Proposal

- Identify all non-redundant phase space functions  $(q, p)$  with classical observables  $\equiv$  *mutables*.
- Evolution of mutables generated by  $\mathcal{H}$ :

$$\dot{q} = \{q, \mathcal{H}\} \quad \dot{p} = \{p, \mathcal{H}\}$$

- Solutions are curves on phase space parametrized by unobservable label  $\tau$ .
- Can use solution to express evolution of any mutable in terms of any other.

$\Rightarrow$  ontological shift: clock variables are considered observable.



# Advantages

- Phase space evolution is integrable in  $\tau$ .
- Can handle “winding numbers”.
- Role of clock is non-ambiguous.



# Quantum Proposal

- Mables become operators that define the quantum observable algebra.

$$(q, p); \{q, p\} = \mathbb{1} \quad \Rightarrow \quad (\hat{q}, \hat{p}); [\hat{q}, \hat{p}]$$

- Evolution of mables generated by  $\hat{\mathcal{H}}$ :

$$\hat{q}(\tau) = e^{-i\hbar\tau\hat{\mathcal{H}}}\hat{q}(0)e^{i\hbar\tau\hat{\mathcal{H}}} \quad \hat{p}(\tau) = e^{-i\hbar\tau\hat{\mathcal{H}}}\hat{p}(0)e^{i\hbar\tau\hat{\mathcal{H}}}$$

$\tau$  - arbitrary unobservable label.

- Equivalent to Schrödinger evolution

$$-i\hbar \frac{\partial \psi}{\partial \tau} = \hat{\mathcal{H}}(\hat{q}, \hat{p})\psi$$

Note:  $\hat{\mathcal{H}}(\hat{q}, \hat{p})$  is  $\tau$ -independent

$\therefore$  “energy” (or some coupling) promoted from coupling constant to conserved quantity.

$\Rightarrow$  classical limit reproduced (energy eigenstate is superselected in classical limit)



# Advantages

- Evolution is unitary in  $\tau$ .
- All potential clocks can be described quantum mechanically.
- Unifying framework for partial/complete observables program:  
⇒ If  $\hat{q}_i = \hat{T}$  is such that  $\hat{T} \approx \langle \hat{T} \rangle \hat{\mathbb{1}}$  then we reproduce standard clock deparametrization
- Our Fundamental Principle is implemented.



# Shape Dynamics

- Split time diffeomorphisms into:
  - ① Hypersurface preserving (i.e., global reparametrizations)
  - ② Hypersurface deforming (i.e., local refoliations)
- Identify:
  - ① Global reparametrization  $\Rightarrow$  evolution
  - ② local refoliations  $\Rightarrow$  redundancy

## Observation

$\exists$  global time such that non-redundant (reduced) phase space is locally scale invariant.

Ontological shift: scale invariance suggests preferred foliation.

$\Rightarrow$  gravity expressed in terms of the dynamics of scale invariant local “shapes”:

One damn shape after another!



# Relational Quantization of Gravity

- Use Shape Dynamics to fix global time.  
⇒ non-redundant phase space is fixed by imposing local scale invariance.
- Relationally quantize Shape Dynamics.





# Summary/Outlook

## Summary

- Fundamental Principle (FP): Time can, in principle, be measured by a clock.  
⇒ Clock values are observable.
- Standard deparametrization in conflict with FP.  
⇒ clocks are not Dirac observables.
- Promote clocks to observables (mutables).
- Mutables evolve unitarily wrt unobservable parameter.
- Can reproduce classical limit, implement FP, and not privilege any clock choice quantum mechanically.

## Outlook

- Mini-superspace quantization.
- Nabu–Goto string quantization: Feynman versus Hadamard propagator.



Thank You!